## Decrease in AI acceptor density in AI-doped 4H-SiC by irradiation with 4.6 MeV electrons

Hideharu Matsuura,<sup>a)</sup> Koichi Aso, Sou Kagamihara, Hirofumi Iwata, and Takuya Ishida Department of Electronic Engineering and Computer Science, Osaka Electro-Communication University, 18-8 Hatsu-cho, Neyagawa, Osaka 572-8530, Japan

Kazuhiro Nishikawa

Sansha Electric Mfg. Co., Ltd., 3-1-56 Nishi-Awaji, Higashi-Yodogawa, Osaka 533-0031, Japan

(Received 14 July 2003; accepted 23 October 2003)

From the temperature dependence of the hole concentration p(T) in a lightly Al-doped 4H-SiC epilayer, an Al acceptor with ~200 meV and an unknown defect with ~370 meV are found. By irradiation with 4.6 MeV electrons, the Al acceptor density is reduced, while the unknown defect density is almost unchanged. This indicates that the substitutional Al atoms in 4H-SiC are displaced by irradiation or that the bonds between the substitutional Al atom and the nearest neighbor atom are broken by irradiation. Moreover, the reduction in p(T) with irradiation arises from the decrease of the Al acceptors, not from the formation of hole traps or donor-like defects. © 2003 American Institute of Physics. [DOI: 10.1063/1.1634381]

Silicon carbide (SiC) is a wide bandgap semiconductor with potential for use in high power and high frequency devices capable of operating at elevated temperatures. Also for electrons with energies greater than 0.5 MeV, the damage constant for lifetime degradation in SiC has been reported to be lower than that in GaAs by more than three orders of magnitude.<sup>1</sup> This indicates a greatly superior resistance of SiC to displacement damage in most radiation environments.

Electron irradiation is an excellent tool for the controlled generation of intrinsic defects in silicon (Si) to be used in high power devices.<sup>2</sup> On the other hand, it degrades the conversion efficiency of Si solar cells used in space.<sup>3–5</sup> While this electron-radiation damage in Si has been investigated by many researchers, the understanding of radiation damage in SiC is far from complete.

The vacancy-related defects induced by electron irradiation were investigated mainly with the following techniques. Using electron paramagnetic resonance spectroscopy (EPR), the following defects induced by electron irradiation were reported;<sup>6–8</sup> a Si Frenkel pair ( $V_{Si}$ -Si), a carbon monovacancy ( $V_C$ ), a divacancy ( $V_C$ - $V_{Si}$ ), and an antisite–vacancy pair ( $C_{Si}$ - $V_C$ ). In electron-irradiated B-doped 6H-SiC, a complex ( $B_{Si}$ - $V_C$ ) of B and a vacancy, which behaved as a deep acceptor,<sup>9</sup> was detected.<sup>10</sup> It was reported that electrons with energies between 1 and 3 MeV were sufficient to displace silicon atoms in 6H-SiC with one electron.<sup>11</sup>

Although the radiation damage has been studied mainly with EPR and deep level transient spectroscopy, the densities and energy levels of hole traps created by irradiation can also be determined accurately from the temperature dependence of the hole concentration p(T).<sup>4</sup> In this article, we report on our investigation of the influence of electron irradiation on p(T) in *p*-type 4H-SiC, since it is important to investigate the cause of radiation degradation of SiC power devices such as metal-oxide-semiconductor field-effect transistors and

bipolar-junction transistors used in space or near nuclear reactors.

A 10  $\mu$ m-thick Al-doped 4H-SiC epilayer (Al-doping density: ~5×10<sup>15</sup> cm<sup>-3</sup>) on *n*-type 4H-SiC (thickness: 375.9  $\mu$ m, resistivity: 0.02  $\Omega$ cm) was cut to a 1×1 cm<sup>2</sup> size. Ohmic metal (Al/Ti) was deposited on the four corners of the surface of the sample, and then the sample was annealed at 900 °C for 1 min in an Ar atmosphere. The p(T) was measured by the van der Pauw method in the temperature range from 135 to 580 K and in a magnetic field of 1.4 T using a modified MMR Technologies' Hall system. After this measurement, the sample was irradiated with 4.6 MeV electrons with 2.6×10<sup>14</sup> cm<sup>-2</sup> fluence at room temperature. Then the p(T) for the sample was measured.

The open circles and diamonds in Fig. 1 represent the measured p(T) for the unirradiated and irradiated samples. It is clear from the figure that the p(T) is reduced by electron irradiation. The possible origins of this reduction in p(T) by irradiation are as follows: (1) a decrease in the acceptor density because substitutional acceptors are moved into the in-



FIG. 1. Temperature dependencies of hole concentration.

4981

Downloaded 22 Dec 2003 to 133.89.3.21. Redistribution subject to AIP license or copyright, see http://ojps.aip.org/aplo/aplcr.jsp

<sup>&</sup>lt;sup>a)</sup>Author to whom correspondence should be addressed; electronic mail: matsuura@isc.osakac.ac.jp

<sup>© 2003</sup> American Institute of Physics



FIG. 2. Temperature dependencies of electron mobility.

terstitial sites or because the bonds between the substitutional acceptor and the nearest neighbor atom are broken, or (2) the creation of hole traps or donor-like defects, which capture holes emitted from the acceptors.

The open circles and diamonds in Fig. 2 represent the experimental temperature dependence of electron mobility  $\mu_n(T)$  for the unirradiated and irradiated samples. The reduction in  $\mu_n(T)$  by irradiation is small above 250 K, but it is large below 250 K. This large reduction below 250 K may result from ion scattering.

In order to investigate acceptor densities and hole-trap densities with p(T), free carrier concentration spectroscopy (FCCS)<sup>5,12–14</sup> is used. The FCCS is a graphical peak analysis method for determining the densities and energy levels of acceptor species and hole traps using p(T) without any assumptions regarding acceptor species and hole traps. Using an experimental p(T), the FCCS signal is defined as<sup>4,5,13,14</sup>

$$H(T, E_{\rm ref}) \equiv \frac{p(T)^2}{(kT)^{5/2}} \exp\left(\frac{E_{\rm ref}}{kT}\right).$$
(1)

The FCCS signal has a peak at the temperature corresponding to each acceptor level or hole-trap level, where  $E_{ref}$  is the parameter that can shift the peak temperature of the FCCS signal within the temperature range of the measurement.

TABLE I. Results obtained by FCCS.

	Before irradiation	After irradiation
$\Delta E_{A1}$ (meV)	203	206
$N_{A1} (\text{cm}^{-3})$	$6.2 \times 10^{15}$	$8.2 \times 10^{14}$
$\Delta E_{A2}$ (meV)	365	383
$N_{A2} ({\rm cm}^{-3})$	$4.2 \times 10^{15}$	$3.4 \times 10^{15}$
$N_{\rm comp}~({\rm cm}^{-3})$	$3.4 \times 10^{13}$	$7.4 \times 10^{14}$

From each peak, the density and energy level of the corresponding acceptor or hole trap can be determined accurately.

The solid line in Fig. 3 represents the FCCS signal of  $H(T, E_{\rm ref})$  with  $E_{\rm ref} = 1.27 \times 10^{-3}$  eV for the unirradiated sample. Here the FCCS signal is calculated by interpolating p(T) with a cubic smoothing natural spline function at intervals of 0.1 K. From the peak at 464.2 K, the corresponding density  $(N_{A2})$  and energy level  $(\Delta E_{A2})$  measured from the valence band  $(E_V)$  are determined as  $4.2 \times 10^{15}$  cm<sup>-3</sup> and 365 meV.

The FCCS signal of  $H2(T, E_{ref})$ , in which the influence of  $\Delta E_{A2}$  is removed, is calculated and denoted by the broken line in Fig. 3. Here,  $E_{ref}$ =0.145 eV. From this peak at 266.7 K, the corresponding density ( $N_{A1}$ ) and energy level ( $\Delta E_{A1}$ ), and the compensating density ( $N_{comp}$ ) are determined as  $6.2 \times 10^{15}$  cm<sup>-3</sup>, 203 meV and  $3.4 \times 10^{13}$  cm<sup>-3</sup>, respectively.

In the same way as illustrated for the unirradiated sample, the densities and energy levels for the irradiated sample are determined and listed in Table I. The value of  $N_{A1}$  is reduced by electron irradiation, while  $N_{A2}$  appears unchanged.

In order to verify the values obtained, the p(T) is simulated with the values shown in Table I. The open circles and diamonds in Fig. 4 represent the experimental p(T) for the unirradiated and irradiated samples, and the solid and broken lines represent the simulated p(T) for the unirradiated and irradiated samples. Each line is in good agreement with the corresponding experimental p(T), indicating that the values determined by FCCS are reliable.

From photoluminescence (PL) measurements<sup>15</sup> and Halleffect measurements,<sup>16</sup> the energy level of ~200 meV is ascribed to the Al acceptor in 4H-SiC. Moreover,  $N_{A1}$  is close



FIG. 3. FCCS signals. FIG. 4. A comparison between experimental and simulated p(T). Downloaded 22 Dec 2003 to 133.89.3.21. Redistribution subject to AIP license or copyright, see http://ojps.aip.org/aplo/aplcr.jsp

to the Al-doping density, suggesting that the origin of the shallower energy level is the Al acceptor. On the other hand, although the possible origin of the deeper energy level is B with which 4H-SiC is sometimes contaminated, the concentration of B in this epilayer, which was determined by secondary ion mass spectroscopy, was  $<4 \times 10^{14}$  cm<sup>-3</sup>, and  $\Delta E_A$  of B was reported to be 285 meV.<sup>16</sup> Although it may be closely linked with the D<sub>1</sub>-line observed by PL,<sup>17</sup> the origin of the deeper energy level is unfortunately unknown to date.

The density of substitutional Al atoms, which act as acceptors, is found to be reduced by 4.6 MeV electron irradiation. This is why the p(T) is decreased by the irradiation. Similar phenomena have been reported in *p*-type Si: the p(T) in B-doped Si was reduced by 10 MeV proton irradiation because of the decrease in the B acceptor density in Si.<sup>5</sup> Moreover, ion implantation has been reported to cause donor-density reduction in nitrogen (N)-doped 4H-SiC, because vacancies created by ion implantation reacted with N donors.<sup>18,19</sup>

The decrease in the Al acceptor density by irradiation is assumed to arise from the movement of the substitutional Al atoms into the interstitial sites, or from the bond-breaking between the substitutional Al atom and the nearest neighbor atom. Further research in this area is in progress.

In summary, the effect of electron irradiation on Aldoped 4H-SiC was investigated with Hall-effect measurements. The hole concentration in *p*-type 4H-SiC was reduced by 4.6 MeV electron irradiation, and by FCCS the  $\sim 200$  meV Al acceptor density was found to be clearly decreased by the irradiation. On the other hand, the unknown  $\sim 370$  meV defect density was unchanged. Therefore, the reduction in p(T) by irradiation resulted from the decrease of Al acceptors, not from the creation of hole traps or donorlike defects.

- <sup>1</sup>A. L. Barry, B. Lehman, D. Fritsch, and D. Bräuning, IEEE Trans. Nucl. Sci. **38**, 1111 (1991).
- <sup>2</sup>S. M. Sze, *Physics of Semiconductor Devices*, 2nd ed. (Wiley, New York, 1981), p. 38 and p. 214.
- <sup>3</sup>T. Hisamatsu, O. Kawasaki, S. Matsuda, T. Nakao, and Y. Wakow, Sol. Energy Mater. Sol. Cells **50**, 331 (1998).
- <sup>4</sup>H. Matsuura, Y. Uchida, T. Hisamatsu, and S. Matsuda, Jpn. J. Appl. Phys., Part 1 37, 6034 (1998).
- <sup>5</sup>H. Matsuura, Y. Uchida, N. Nagai, T. Hisamatsu, T. Aburaya, and S. Matsuda, Appl. Phys. Lett. **76**, 2092 (2000).
- <sup>6</sup>H. J. von Bardeleben and J. L. Cantin, Mater. Sci. Forum **353–356**, 513 (2001).
- <sup>7</sup>H. J. von Bardeleben, J. L. Cantin, P. Baranov, and E. N. Mokhov, Mater. Sci. Forum **353–356**, 509 (2001).
- <sup>8</sup>Th. Lingner, S. Greulich-Weber, and J.-M. Spaeth, Mater. Sci. Forum **353–356**, 505 (2001).
- <sup>9</sup>I. V. Ilyin, E. N. Mokhov, and P. G. Baranov, Mater. Sci. Forum **353–356**, 521 (2001).
- <sup>10</sup> V. Ya. Bratus, I. N. Makeeva, S. M. Okulov, T. L. Petrenko, T. T. Petrenko, and H. J. von Bardeleben, Mater. Sci. Forum **353–356**, 517 (2001).
- <sup>11</sup>A. A. Rempel, K. Blaurock, K. J. Reichle, W. Sprengel, and H.-E. Schaefer, Mater. Sci. Forum **389–393**, 485 (2002).
- <sup>12</sup>H. Matsuura and K. Sonoi, Jpn. J. Appl. Phys., Part 2 35, L555 (1996).
- <sup>13</sup> H. Matsuura, Y. Masuda, Yi Chen, and S. Nishino, Jpn. J. Appl. Phys., Part 1 **39**, 5069 (2000).
- <sup>14</sup>H. Matsuura, K. Morita, K. Nishikawa, T. Mizukoshi, M. Segawa, and W.
- Susaki, Jpn. J. Appl. Phys., Part 1 **41**, 496 (2002). <sup>15</sup> M. Ikeda, H. Matsunami, and T. Tanaka, Phys. Rev. B **22**, 2842 (1980).
- <sup>16</sup>T. Troffer, M. Schadt, T. Frank, H. Itoh, G. Pensl, J. Heindl, H. P. Strunk, and M. Maier, Phys. Status Solidi A **162**, 277 (1997).
- <sup>17</sup>T. Egilsson, J. P. Bergman, I. G. Ivanov, A. Henry, and E. Janzén, Phys. Rev. B **59**, 1956 (1999).
- <sup>18</sup> A. Hallén, A. Henry, P. Pellegrino, B. G. Svensson, and D. Aberg, Mater. Sci. Eng., B **61–62**, 378 (1999).
- <sup>19</sup>D. Åberg, A. Hallén, P. Pellegrino, and G. Svensson, Appl. Phys. Lett. 78, 2908 (2001).