

# Background

In order to use SiC wafers or epilayers for electronic devices, **an accurate evaluation of densities and energy levels of electrically active impurities and defects is essential.**

## Electrical method for characterizations

For **deep impurities or defects**

- a. deep level transient spectroscopy (**DLTS**)
- b. isothermal capacitance transient spectroscopy (**ICTS**)

For **shallow impurities or defects** ??

**The temperature dependence of the majority-carrier concentration  $n(T)$  includes a lot of important information on shallow impurities or defects.**

**Curve-fitting method  $\longrightarrow$  unique solution is not obtained**

## Aim

1. To know **the number of types** of shallow impurities or defects in SiC.
2. To determine **the densities and energy levels** of these impurities or defects.

DLTS or ICTS can **uniquely** determine the densities and energy levels of **deep impurities or defects**.



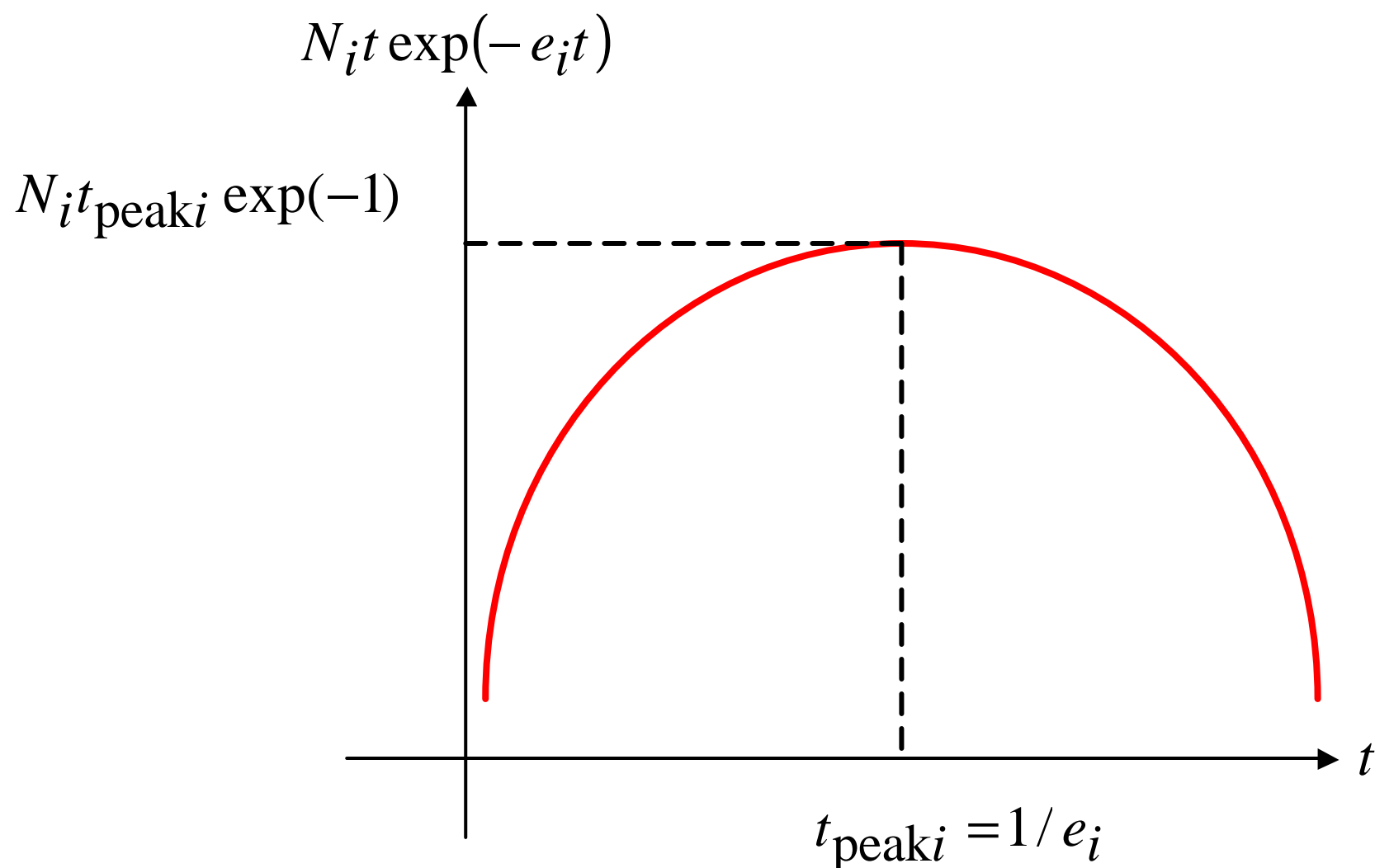
This is because **each peak** in the signal corresponds to a **one-to-one impurity or defect**.

In ICTS, the signal is defined as

$$S(t) \equiv t \frac{dC(t)}{dt}. \quad C(t): \text{transient capacitance}$$

The signal can be theoretically expressed as

$$S(t) \propto \sum_i N_i t \exp(-e_i t). \quad \begin{array}{l} N_i : \text{density of } i\text{-th trap} \\ e_i : \text{emission rate of } i\text{-th trap} \end{array}$$

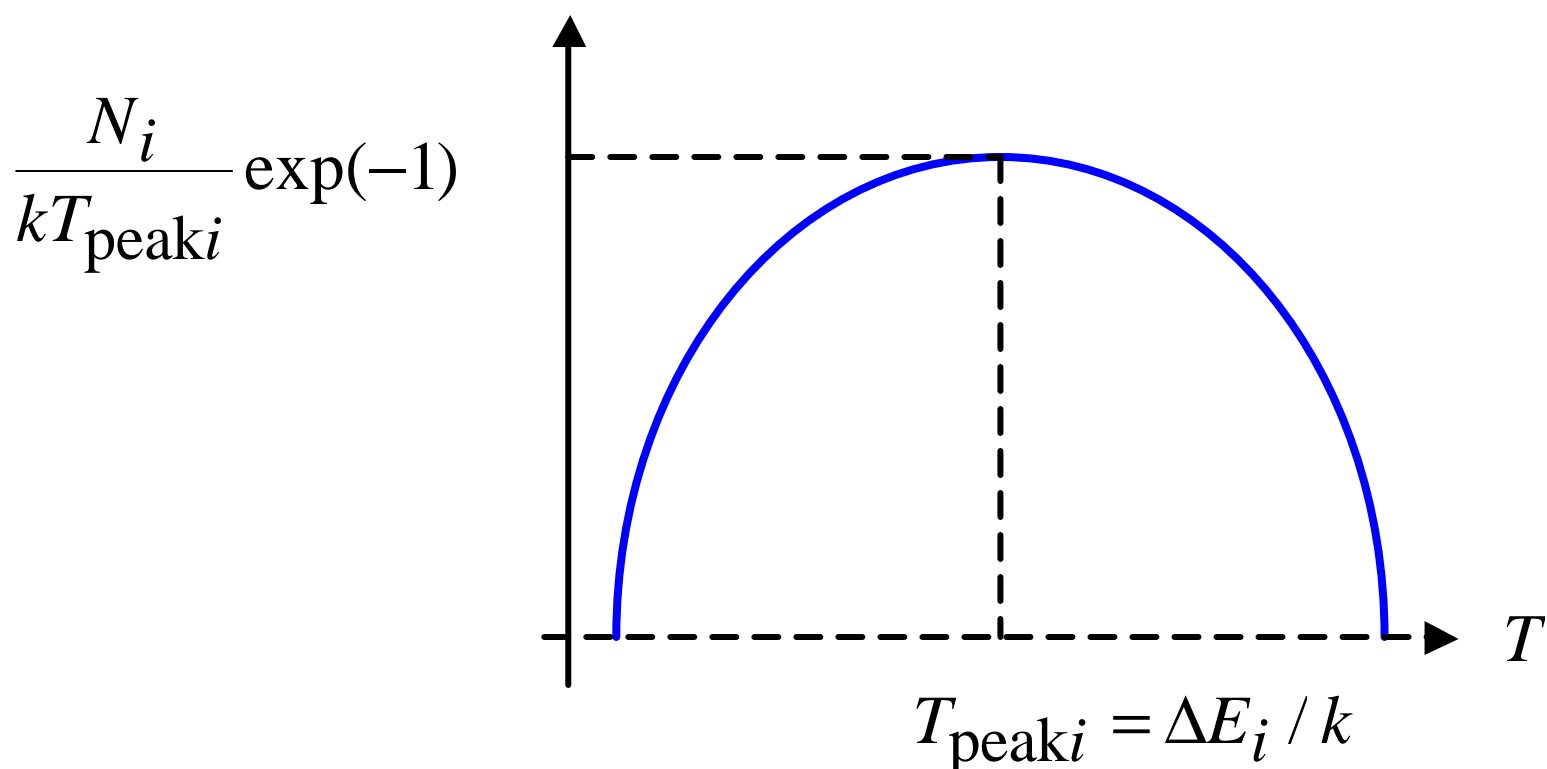


$N_i t \exp(-e_i t)$  plays an **important role** in the ICTS analysis.

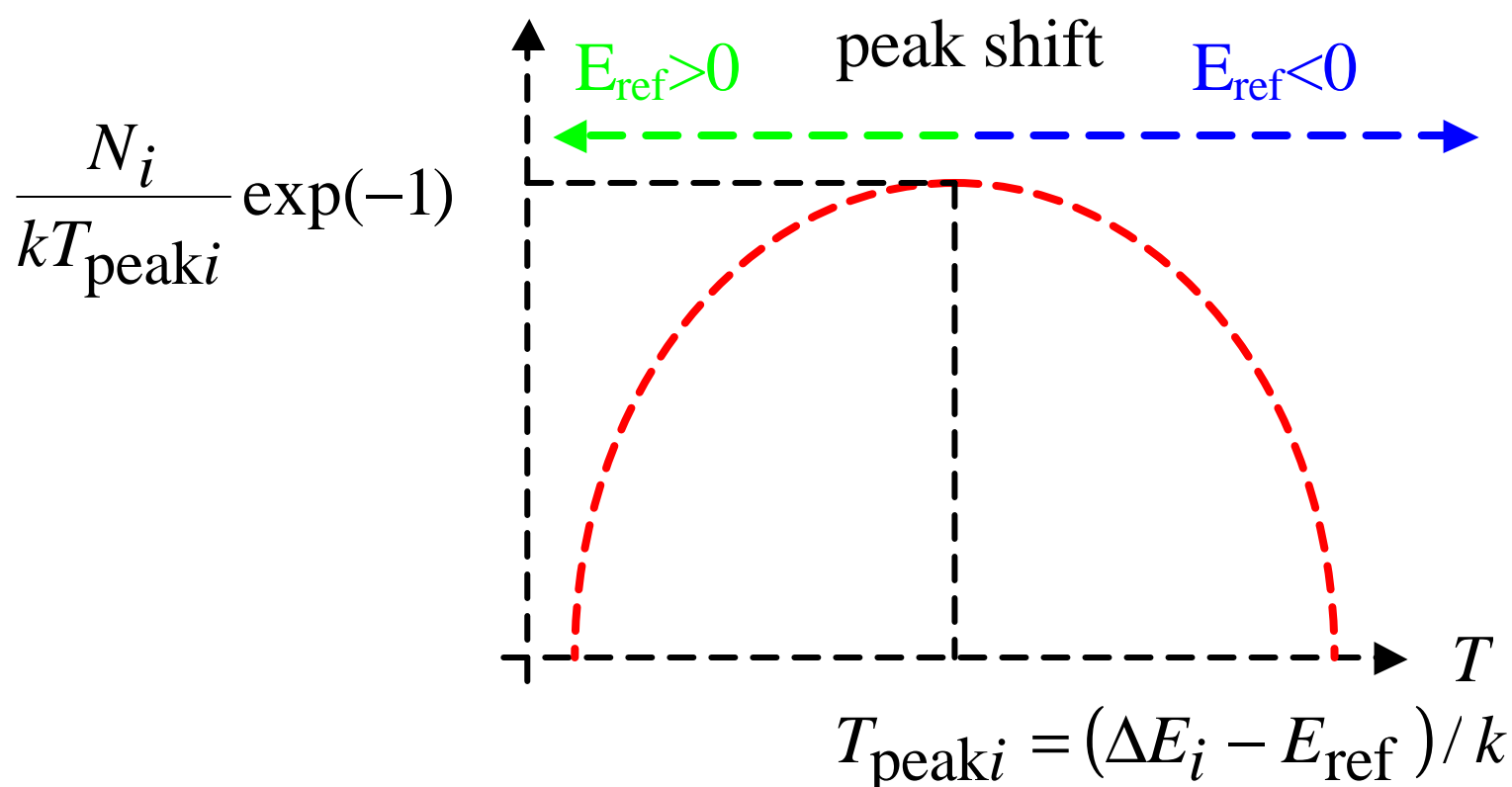
# Concept of New Graphical Method

Let us consider the following functions:

$$1. \frac{N_i}{kT} \exp\left(-\frac{\Delta E_i}{kT}\right)$$



$$2. \frac{N_i}{kT} \exp\left(-\frac{\Delta E_i - E_{\text{ref}}}{kT}\right)$$



Introducing a **parameter**  $E_{\text{ref}}$ , we can **change the peak condition**.

# Desirable Function to be evaluated

The temperature dependence of the electron concentration  $n(T)$  in an n-type semiconductor is described as follows:

1. From the charge neutrality condition,

$$n(T) = \sum_i N_{Di} [1 - f(\Delta E_{Di})] - N_A \quad (1)$$

$N_{Di}$ :  $i$ -th donor density

$\Delta E_{Di}$ :  $i$ -th donor level measured from  $E_C$

$f(\Delta E_{Di})$ : Fermi-Dirac distribution function

$N_A$ : acceptor density

2. From the effective density of states  $N_C(T)$ ,

$$n(T) = N_C(T) \exp\left(-\frac{\Delta E_F}{kT}\right) \quad (2)$$

$\Delta E_F$ : Fermi level measured from  $E_C$

Using the two equations (1) and (2),  
how shall we define a desirable function?

↓  
[ The defined function should include the following terms. ]

$$\sum_i \frac{N_{Di}}{kT} \exp\left(-\frac{\Delta E_{Di} - E_{\text{ref}}}{kT}\right)$$

## Desirable Definition

When we define a function to be evaluated as

$$H(T, E_{\text{ref}}) \equiv \frac{n(T)^2}{(kT)^{2.5}} \exp\left(\frac{E_{\text{ref}}}{kT}\right),$$

the defined function can be expressed as

$$H(T, E_{\text{ref}}) = \sum_i \frac{N_{Di}}{kT} \exp\left(-\frac{\Delta E_{Di} - E_{\text{ref}}}{kT}\right) I(E_{Di})$$

$$- N_A \frac{N_{C0}}{kT} \exp\left(\frac{E_{\text{ref}} - \Delta E_F}{kT}\right)$$

using eqs. (1) and (2), where  $N_{C0} = 2(2\pi m_n^* / h^2)^{1.5}$ .

Since  $I(\Delta E_{Di})$  is found to be almost independent of  $T$ ,

$H(T, E_{\text{ref}})$  can be approximately expressed as

$$\text{the sum of } \frac{N_{Di}}{kT} \exp\left(-\frac{\Delta E_{Di} - E_{\text{ref}}}{kT}\right).$$

## Good points of our method

1. Definition: 
$$H(T, E_{\text{ref}}) \equiv \frac{n(T)^2}{(kT)^{2.5}} \exp\left(\frac{E_{\text{ref}}}{kT}\right)$$

$H(T, E_{\text{ref}})$  has peaks corresponding to each energy level.

$i$ -th peak temperature  $\longrightarrow$  energy level of  $i$ -th donor

$i$ -th peak value  $\longrightarrow$  density of  $i$ -th donor

2. Peak temperature: 
$$T_{\text{peak}i} \cong \frac{\Delta E_{\text{Di}} - E_{\text{ref}}}{k}$$

Even when none of peaks of  $H(T, 0)$  appear within the measurement temperature range, the peak of  $H(T, E_{\text{ref}})$  can be shifted to the measurement temperature range by changing  $E_{\text{ref}}$ .

3. From the number of peaks in  $H(T, E_{\text{ref}})$ , we can know the number of types of donors.

4. We can determine the compensating density (i.e., acceptor density).

# Undoped 3C-SiC

## Growth condition

(Atmospheric pressure chemical vapor deposition)

### 1. Cleaning of (100) n-type **Si substrate**

(Etching of Si surface)

1175 °C, 11 min.

HCl: 63 sccm, H<sub>2</sub>: 1.5 slm

### 2. Formation of buffer layer on Si substrate

(Carbonization of Si surface)

1350 °C, 3 min.

C<sub>3</sub>H<sub>8</sub>: 1 sccm, H<sub>2</sub>: 1 slm

### 3. Growth of **undoped 3C-SiC**

1350 °C

**Si<sub>2</sub>(CH<sub>3</sub>)<sub>6</sub>**: 0.5 sccm, H<sub>2</sub>: 2.5 slm

growth rate: 4.3 μm/h

## Conditions of Hall-effect measurement

**Removal of Si substrate** (chemical etching)

Thicknesses: 8 μm, 16 μm, 32 μm

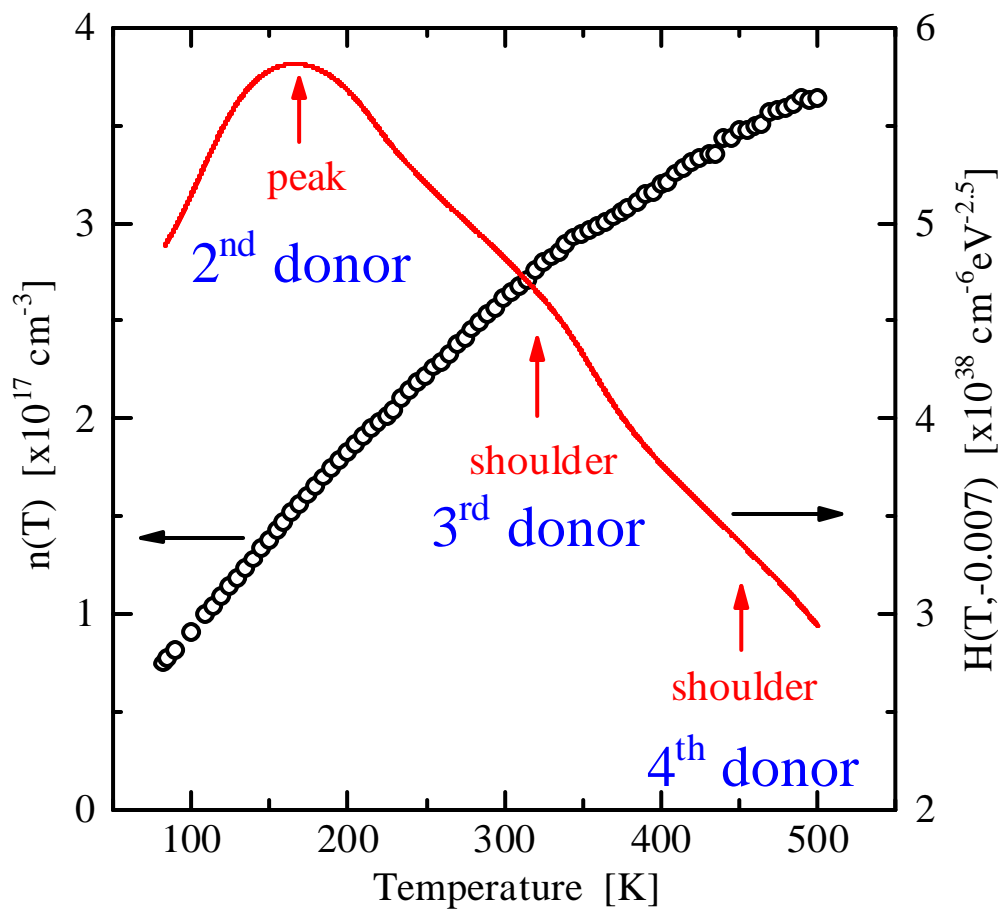
Size: 5 × 5 mm<sup>2</sup>

Magnetic field: 5 kG

Temperature range: 85 K~500 K

# Electron Concentration and Fermi Level

$$H(T, E_{\text{ref}}) \equiv \frac{n(T)^2}{(kT)^{2.5}} \exp\left(\frac{E_{\text{ref}}}{kT}\right)$$



peak

166 K

$5.8 \times 10^{38} \text{ cm}^{-6} \text{ eV}^{-2.5}$

2<sup>nd</sup> donor level

$\Delta E_{D2} = 46 \text{ meV}$

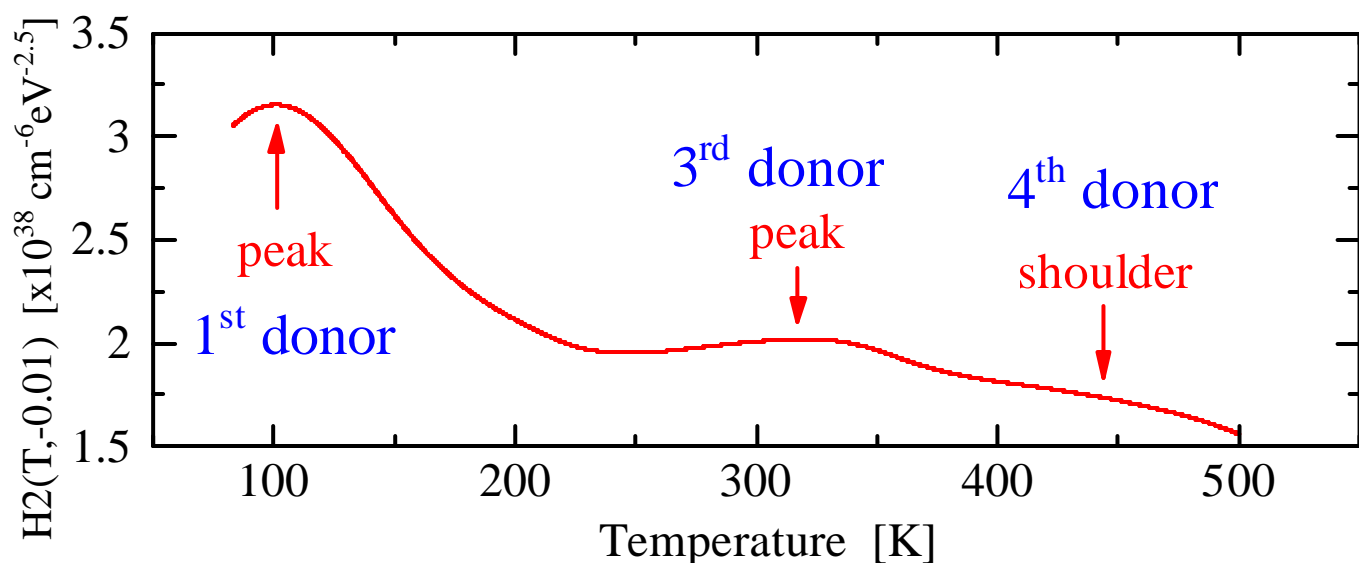
2<sup>nd</sup> donor density

$N_{D2} = 1.7 \times 10^{17} \text{ cm}^{-3}$

$H_2(T, E_{\text{ref}})$  that is not influenced by the 2<sup>nd</sup> donor:

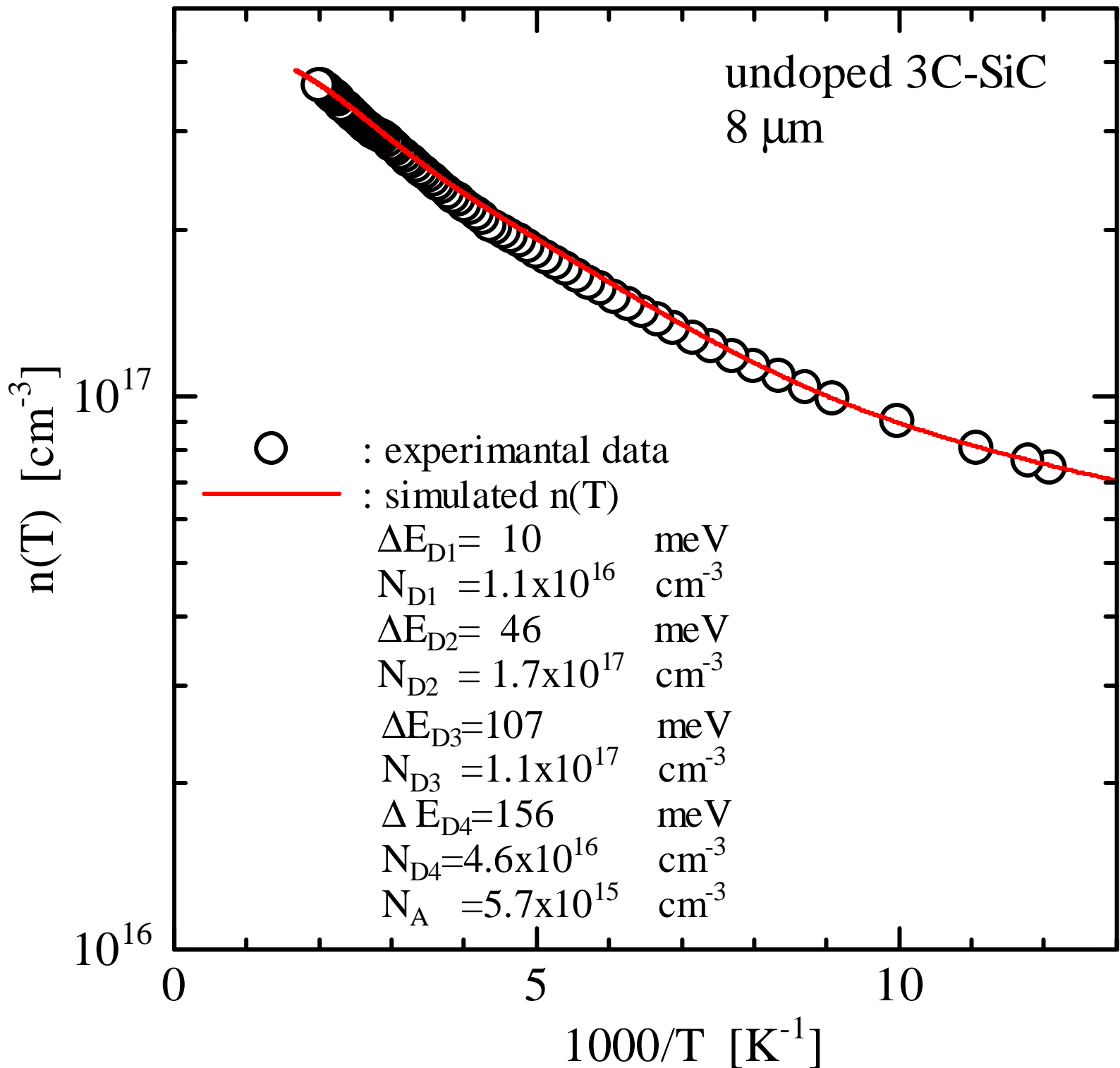
$$H_2(T, E_{\text{ref}}) \equiv \frac{n(T)^2}{(kT)^{2.5}} \exp\left(\frac{E_{\text{ref}}}{kT}\right)$$

$$- \frac{N_{D2}}{kT} \exp\left(-\frac{\Delta E_{D2} - E_{\text{ref}}}{kT}\right) I(\Delta E_{D2})$$





**Comparison of the experimental  $n(T)$  with the  $n(T)$  simulated with the values determined by  $H(T, E_{\text{ref}})$**



The simulated  $n(T)$  is quantitatively **in good agreement with** the experimental  $n(T)$ .



The values determined by  $H(T, E_{\text{ref}})$  are **reasonable**.

## Thickness Dependence of Donor Densities and Donor Levels

3C-SiC thickness [ $\mu\text{m}$ ]	8	16	32
$E_{D1}$ [meV]	10	7	14
$N_{D1}$ [ $\times 10^{16} \text{cm}^{-3}$ ]	11	8.1	4.7
$E_{D2}$ [meV]	46	46	54
$N_{D2}$ [ $\times 10^{16} \text{cm}^{-3}$ ]	17	20	8.1
$E_{D3}$ [meV]	107	97	120
$N_{D3}$ [ $\times 10^{16} \text{cm}^{-3}$ ]	11	13	10
$E_{D4}$ [meV]	156	-----	-----
$N_{D4}$ [ $\times 10^{16} \text{cm}^{-3}$ ]	4.6	-----	-----
$N_A$ [ $\times 10^{16} \text{cm}^{-3}$ ]	1.3	0.99	0.57

### Origin of donors

#### 7-14 meV donor

**defect-impurity complex or nonstoichiometric defect ?**

(this donor was reported from Hall-effect measurements in undoped 3C-SiC grown from a mixture of  $\text{SiH}_4$  and  $\text{C}_3\text{H}_8$ , where the donor density was higher than  $10^{18} \text{cm}^{-3}$  and the compensation ratio was higher than 0.9)

#### 46-54 meV donor

**substitutional nitrogen atom**

(this donor was reported from photoluminescence)

#### 97-120 meV donor

#### 156 meV donor

}  $\longrightarrow$  not reported yet

## Conclusions

Using  $n(T)$  obtained by Hall-effect measurements, we have attempted to precisely determine the densities and energy levels of shallow electrically active impurities and defects.

When we define a function to be evaluated as

$$H(T, E_{\text{ref}}) \equiv \frac{n(T)^2}{(kT)^{2.5}} \exp\left(\frac{E_{\text{ref}}}{kT}\right),$$

using peaks of  $H(T, E_{\text{ref}})$

1. we can determine **the number of types** of shallow impurities or defects in a semiconductor,
2. we can precisely determine **the density and energy level** of each impurity or defect,
3. we can determine **the compensating density**, and
4. we can easily **verify** the obtained values.

## References

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- 3) H. Matsuura, Y. Uchida, T. Hisamatsu and S. Matsuda: Jpn. J. Appl. Phys. **37** (1998) 6034.
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- 6) H. Matsuura, Y. Uchida, N. Nagai, T. Hisamatsu, T. Aburaya and S. Matsuda: Appl. Phys. Lett. **76** (2000) 2092.
- 7) H. Matsuura, Y. Masuda, Y. Chen and S. Nishino: Jpn. J. Appl. Phys. **39** (2000) to be published in September.

More information is in our web site (<http://www.osakac.ac.jp/labs/matsuura>).

Free application software of this method for Windows OS is to appear  
in our web site by the end of this year.